

Any notes/figures demonstrated below are self-contained and included here for the purpose of completeness - some of which may be absent from the main lecture materials. Please review with discretion.

## S1.1 Proofs and Verifications of Some Properties Seen in Tensor Analysis

### 1. Proof of $\varepsilon_{ijk}\varepsilon_{ipq} = \delta_{jp}\delta_{kq} - \delta_{kp}\delta_{jq}$

From the lecture, you have seen the contraction property of two Levi-Civita symbols sharing one same index is  $\varepsilon_{ijk}\varepsilon_{ipq} = \delta_{jp}\delta_{kq} - \delta_{kp}\delta_{jq}$ . The following section shall explore the derivation of this property.

We first introduce three orthonormal bases

$$\mathbf{e}_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad \mathbf{e}_2 = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, \quad \mathbf{e}_3 = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}.$$

With the three orthogonal bases defined above, we can re-write the Kronecker delta as

$$\delta_{ij} = \mathbf{e}_i^T \mathbf{e}_j,$$

where  $i$  and  $j$  are two free indices ranging from 1 to 3.

#### Examples to verify $\delta_{ij} = \mathbf{e}_i^T \mathbf{e}_j$

- For the case  $i = 1, j = 1$ ,

$$\mathbf{e}_1^T \mathbf{e}_1 = \mathbf{e}_1^T \mathbf{e}_2 = \begin{pmatrix} 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = 1 \quad \equiv \quad \delta_{11} = 1 \quad \checkmark$$

- For the case  $i = 1, j = 2$ ,

$$\mathbf{e}_1^T \mathbf{e}_2 = \mathbf{e}_1^T \mathbf{e}_3 = \begin{pmatrix} 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = 0 \quad \equiv \quad \delta_{12} = 0 \quad \checkmark$$

Same for other combinations of  $i$  and  $j$ .

Similarly, we can re-write the Levi-Civita symbol as

$$\varepsilon_{ijk} = |\mathbf{e}_i \ \mathbf{e}_j \ \mathbf{e}_k|,$$

where  $|\star|$  denotes the determinant of the matrix.

#### Examples to verify $\varepsilon_{ijk} = |\mathbf{e}_i \ \mathbf{e}_j \ \mathbf{e}_k|$

- For the case  $i = 1, j = 1, k = 3$ ,

$$|\mathbf{e}_1 \ \mathbf{e}_j \ \mathbf{e}_k| = |\mathbf{e}_1 \ \mathbf{e}_1 \ \mathbf{e}_3| = \begin{vmatrix} 1 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{vmatrix} = 0 \quad \equiv \quad \varepsilon_{113} = 0 \quad \checkmark$$

- For the case  $i = 1, j = 2, k = 3$ ,

$$|\mathbf{e}_1 \ \mathbf{e}_j \ \mathbf{e}_k| = |\mathbf{e}_1 \ \mathbf{e}_2 \ \mathbf{e}_3| = \begin{vmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{vmatrix} = 1 \quad \equiv \quad \varepsilon_{123} = 1 \quad \checkmark$$

Same for other combinations of  $i, j, k$ .

With the identities above, the product of two Levi-Civitas (sharing the same 'i' index) are

$$\varepsilon_{ijk}\varepsilon_{ipq} = |\mathbf{e}_i \ \mathbf{e}_j \ \mathbf{e}_k| \begin{pmatrix} \mathbf{e}_i^T \\ \mathbf{e}_p^T \\ \mathbf{e}_q^T \end{pmatrix} = \begin{vmatrix} \mathbf{e}_i^T \mathbf{e}_i & \mathbf{e}_i^T \mathbf{e}_j & \mathbf{e}_i^T \mathbf{e}_k \\ \mathbf{e}_p^T \mathbf{e}_i & \mathbf{e}_p^T \mathbf{e}_j & \mathbf{e}_p^T \mathbf{e}_k \\ \mathbf{e}_q^T \mathbf{e}_i & \mathbf{e}_q^T \mathbf{e}_j & \mathbf{e}_q^T \mathbf{e}_k \end{vmatrix} = \begin{vmatrix} \delta_{ii} & \delta_{ij} & \delta_{ik} \\ \delta_{pi} & \delta_{pj} & \delta_{pk} \\ \delta_{qi} & \delta_{qj} & \delta_{qk} \end{vmatrix}.$$

For the case  $i \neq j \neq k$ , we know  $\delta_{ii} = 1$ ,  $\delta_{ij} = 0$ ,  $\delta_{ik} = 0$ , leading the determinant simplified as

$$\epsilon_{ijk}\epsilon_{ipq} = \begin{vmatrix} 1 & 0 & 0 \\ \delta_{pi} & \delta_{pj} & \delta_{pk} \\ \delta_{qi} & \delta_{qj} & \delta_{qk} \end{vmatrix} = 1 \cdot (\delta_{pj}\delta_{qk} - \delta_{pk}\delta_{jq}),$$

which proves the contraction property as desired.

## 2. Verify $(\mathbf{a} \times \mathbf{b})_i = \epsilon_{ijk}a_jb_k$ by exhaustion

Let the vector  $\mathbf{a} = \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix}$  and  $\mathbf{b} = \begin{pmatrix} b_1 \\ b_2 \\ b_3 \end{pmatrix}$ . By the right hand rule, the product of  $\mathbf{a}$  and  $\mathbf{b}$  is  $\mathbf{a} \times \mathbf{b} = \begin{pmatrix} a_2b_3 - a_3b_2 \\ a_3b_1 - a_1b_3 \\ a_1b_2 - a_2b_1 \end{pmatrix}$ .

From the lecture, you have seen  $\mathbf{a} \times \mathbf{b}$  can also be expressed in tensor notation:  $(\mathbf{a} \times \mathbf{b})_i = \epsilon_{ijk}a_jb_k$ . The following section shall explore why this formula works by exhausting the possible values of indices.

Each index  $i, j, k$  can freely range from 1 to 3. For the case  $i = 1$ , by  $(\mathbf{a} \times \mathbf{b})_{i=1} = \epsilon_{1jk}a_jb_k$ , we can list all combinations of  $j = \{1, 2, 3\}$  and  $k = \{1, 2, 3\}$

	$j = 1$	$j = 2$	$j = 3$
$k = 1$	$\epsilon_{111}a_1b_1$	$\epsilon_{121}a_2b_1$	$\epsilon_{131}a_3b_1$
$k = 2$	$\epsilon_{112}a_1b_2$	$\epsilon_{122}a_2b_2$	$\epsilon_{132}a_3b_2$
$k = 3$	$\epsilon_{113}a_1b_3$	$\epsilon_{123}a_2b_3$	$\epsilon_{133}a_3b_3$

Recall the definition of the Levi-Civita symbol:  $\epsilon_{ijk} = \begin{cases} 1, & \{i, j, k\} = \{1, 2, 3\}, \{2, 3, 1\}, \{3, 1, 2\} \\ -1, & \{i, j, k\} = \{3, 2, 1\}, \{2, 1, 3\}, \{1, 3, 2\} \\ 0, & \text{otherwise} \end{cases}$ . Hence, the only two non-zero terms left in the table above are  $\epsilon_{132} = -1$  and  $\epsilon_{123} = 1$ ,

	$j = 1$	$j = 2$	$j = 3$
$k = 1$	0	0	0
$k = 2$	0	0	$-a_3b_2$
$k = 3$	0	$a_2b_3$	0

Summing these two terms:  $(\mathbf{a} \times \mathbf{b})_{i=1} = a_2b_3 - a_3b_2$ , **which is the 1<sup>st</sup> component of the  $\mathbf{a} \times \mathbf{b}$  vector.**

Following the same procedure to exhaust the combinations of  $j$  and  $k$  under  $i = 2$  and  $i = 3$  surely yields the result that 'coincides' with the  $i^{\text{th}}$  vector components in  $\mathbf{a} \times \mathbf{b}$ .

## 3. Proof of $\nabla \times \nabla\phi = 0$

By writing  $\nabla \times \nabla\phi$  in index notation,

$$\epsilon_{ijk} \frac{\partial}{\partial x_j} (\nabla\phi)_k = \epsilon_{ijk} \frac{\partial^2 \phi}{\partial x_j \partial x_k} \quad (\text{A})$$

Swipe the notation in equation A:  $j \leftrightarrow k$

$$\epsilon_{ijk} \frac{\partial^2 \phi}{\partial x_j \partial x_k} = \epsilon_{ikj} \frac{\partial^2 \phi}{\partial x_k \partial x_j} \quad (\text{B})$$

Also, by the commutative property of partial differentiation, equation A can be written as

$$\epsilon_{ijk} \frac{\partial^2 \phi}{\partial x_j \partial x_k} = \epsilon_{ijk} \frac{\partial^2 \phi}{\partial x_k \partial x_j} \quad (\text{C})$$

Equate equation B and C:  $\epsilon_{ikj} \frac{\partial^2 \phi}{\partial x_k \partial x_j} = \epsilon_{ijk} \frac{\partial^2 \phi}{\partial x_k \partial x_j}$ . Since  $\epsilon_{ijk} = -\epsilon_{ikj}$ , the only circumstance to equate the L.H.S. and R.H.S. is both terms are 0. (cf.  $-a = a$  indicates  $a = 0$ ). Hence,  $\nabla \times \nabla\phi = 0$ .